

Automorphic string amplitudes

Henrik Gustafsson

String Theory Seminar
Oxford, Jan 2017

 hgustafsson.se

Based on

Small automorphic representations and degenerate Whittaker vectors

HG, Axel Kleinschmidt, Daniel Persson

[arXiv:1412.5625](https://arxiv.org/abs/1412.5625) [math.NT]

[GKP14]

Journal of Number Theory 166 (Sep, 2016) 344–399

Eisenstein series and automorphic representations

Philipp Fleig, HG, Axel Kleinschmidt, Daniel Persson

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Cambridge University Press (2017)

Upcoming work with

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E_6, E_7, E_8

Outline

Compute Fourier coefficients of automorphic forms to capture information about non-perturbative effects such as instantons and black holes

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- Scattering amplitudes
4-graviton | Derivative expansion | U-duality | SUSY constraints

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Motivation

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- Hecke eigenvalues
- Point counts of elliptic curves
- Langlands program
L-functions | The Langlands–Shahidi method

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Scattering amplitudes | Black hole microstate counting
- Statistical mechanics
Two-dimensional models of crystals

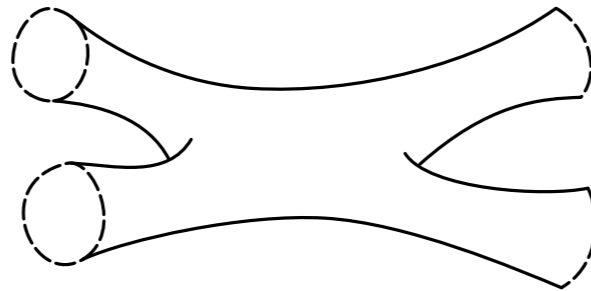
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String theory

Toroidal compactifications of type IIB string theory

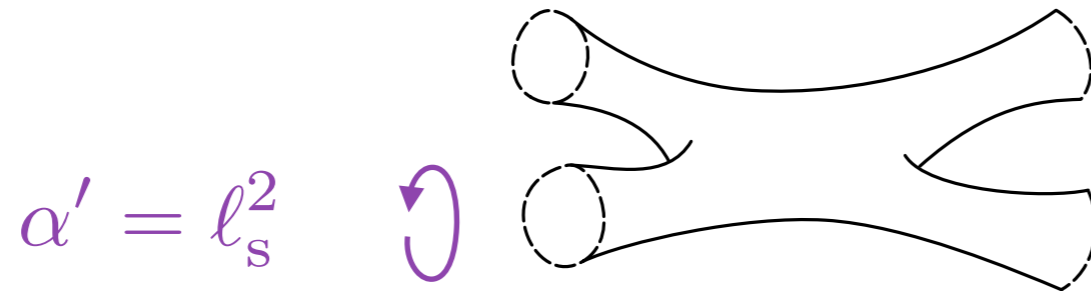
4-graviton scattering amplitudes



String theory

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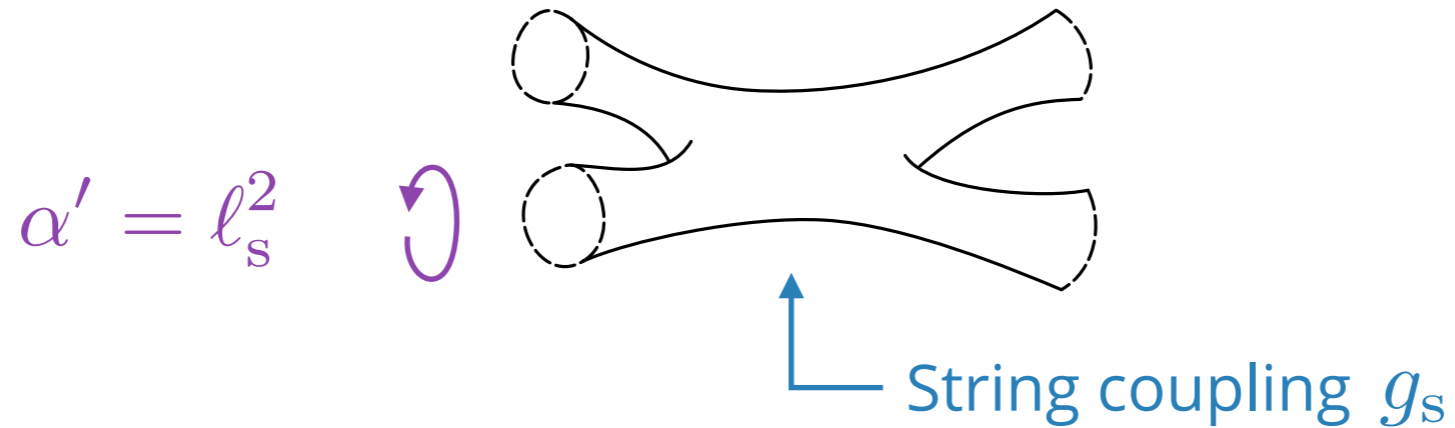
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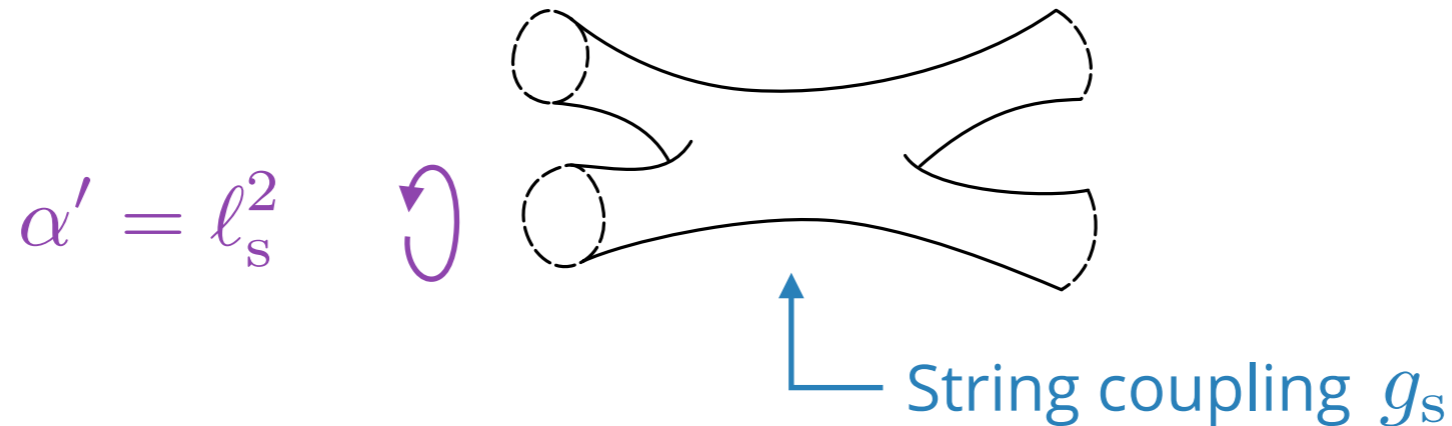
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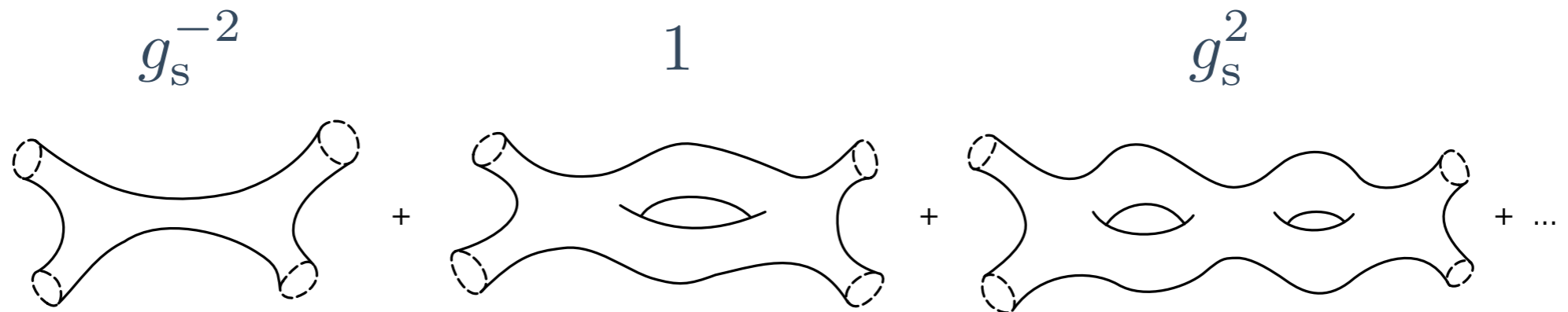


$$s = -\frac{\alpha'}{4}(k_1 + k_2)^2$$

$$t = -\frac{\alpha'}{4}(k_1 + k_3)^2$$

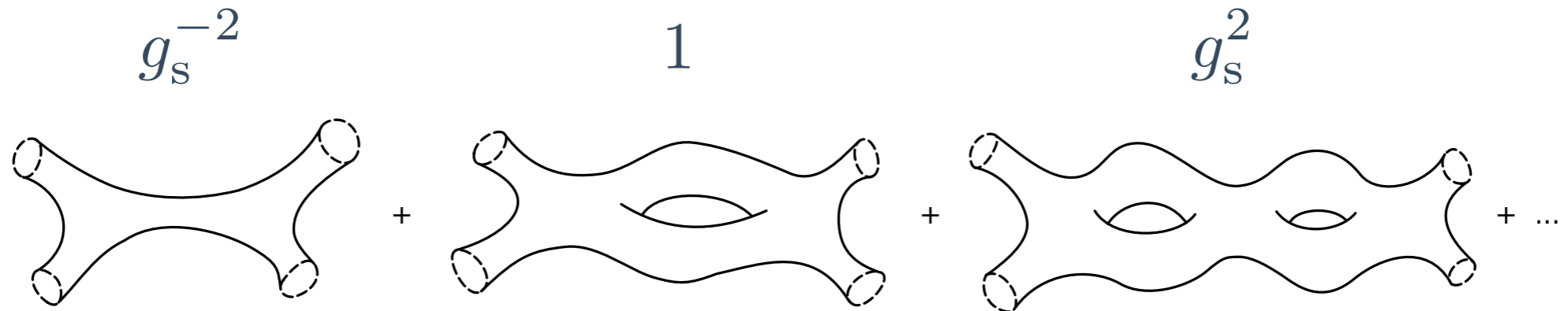
$$u = -\frac{\alpha'}{4}(k_1 + k_4)^2$$

Interactions



4-graviton amplitude in 10 dimensions:

Interactions

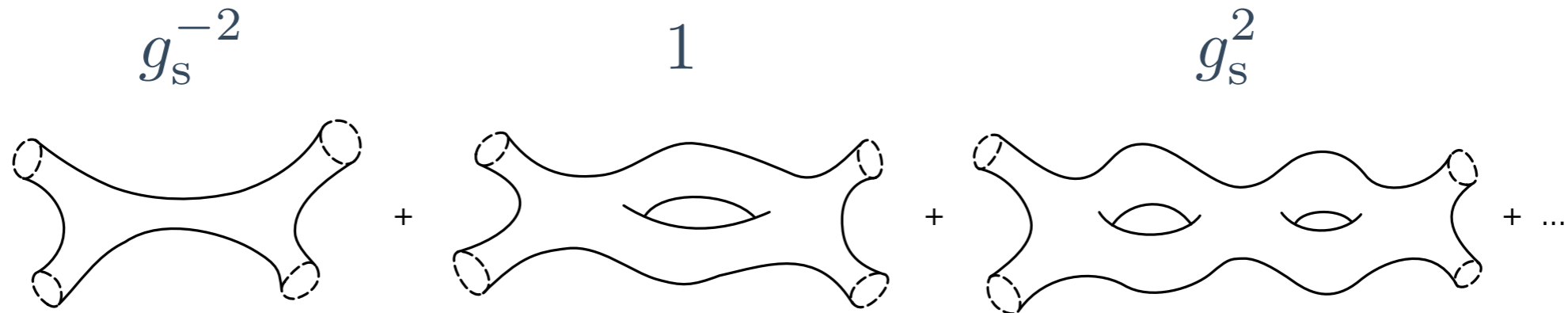


4-graviton amplitude in 10 dimensions:

$$\mathcal{A} = \left(g_s^{-2} \frac{1}{stu} \frac{\Gamma(1-s)\Gamma(1-t)\Gamma(1-u)}{\Gamma(1+s)\Gamma(1+t)\Gamma(1+u)} \right) \mathcal{R}^4$$

[Green-Schwarz, Green-Schwarz-Brink, Gross-Witten]

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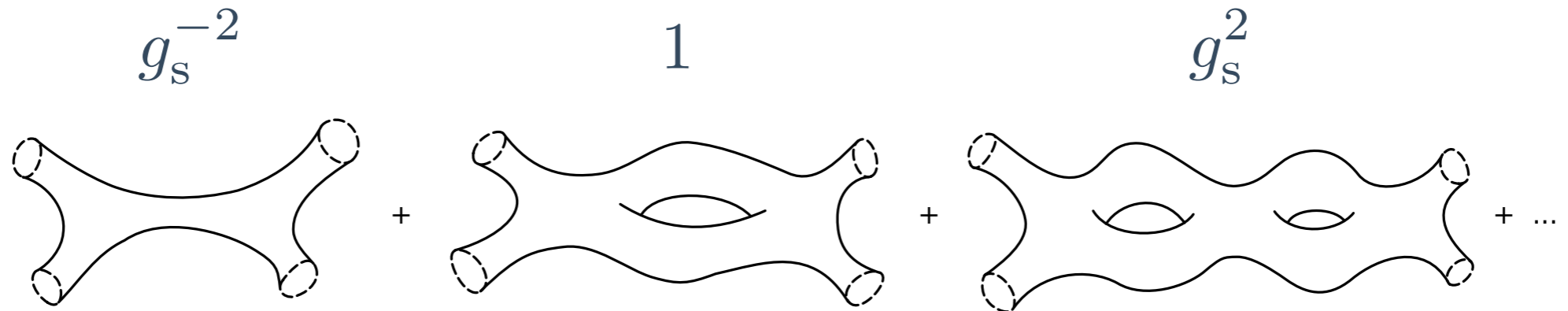
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↗ Contraction of 4 linearized
 Riemann tensors and
 standard rank 8 tensors
 $t_8 t_8 \mathcal{R}^4$

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Interactions

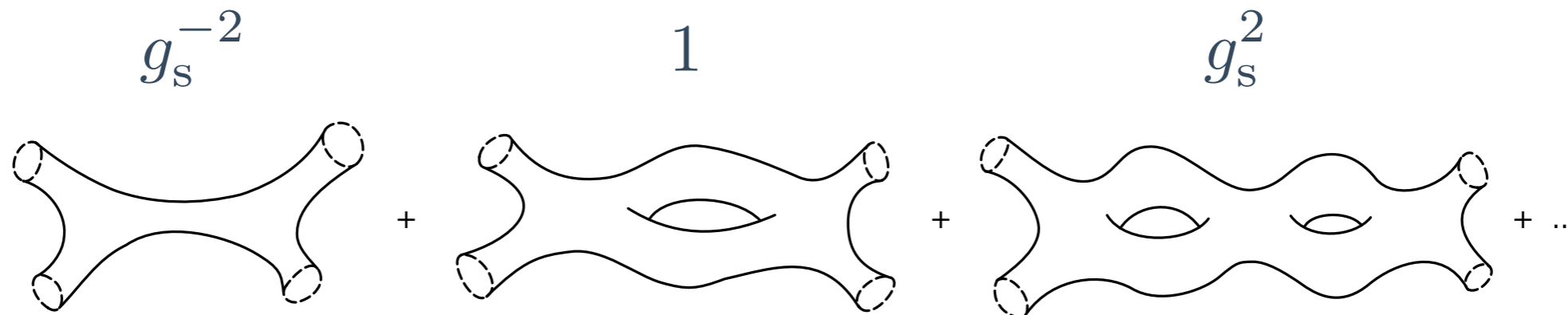


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$$\mathcal{A} = \left(g_s^{-2} \frac{1}{stu} \frac{\Gamma(1-s)\Gamma(1-t)\Gamma(1-u)}{\Gamma(1+s)\Gamma(1+t)\Gamma(1+u)} + 2\pi \int_{\mathcal{F}} \frac{d^2\tau}{(\text{Im } \tau)^2} \mathcal{B}_1(s, t, u; \tau) \right) \mathcal{R}^4$$

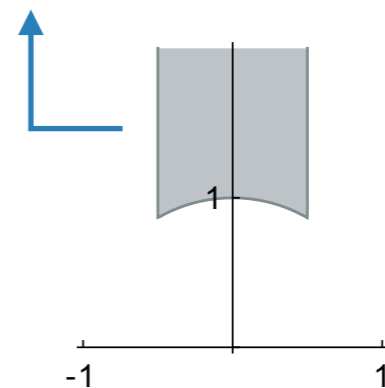
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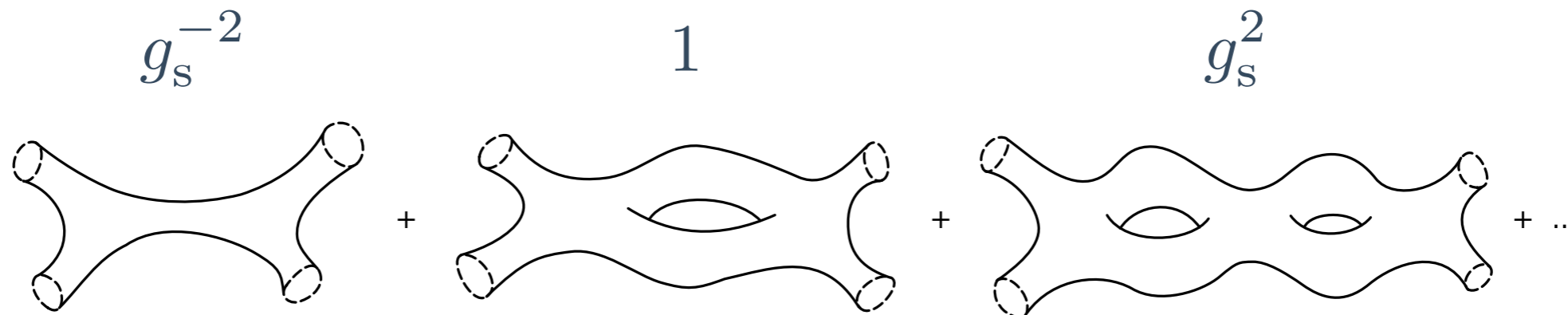
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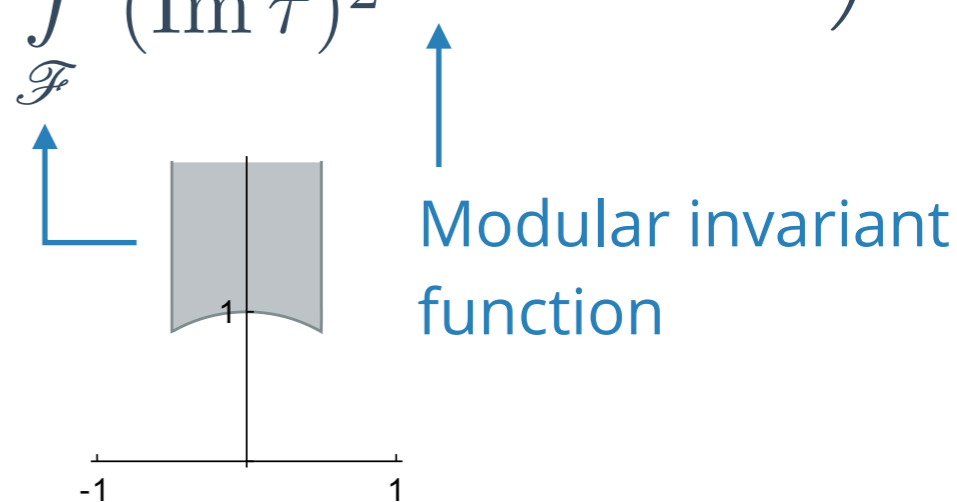
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4-graviton amplitude in 10 dimensions:

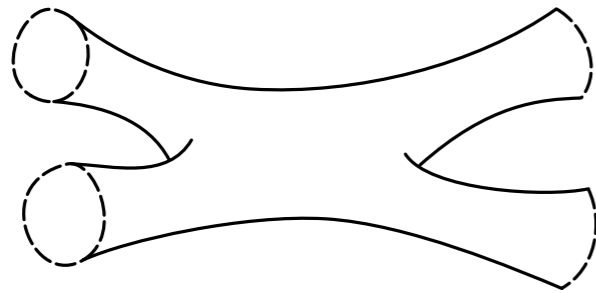
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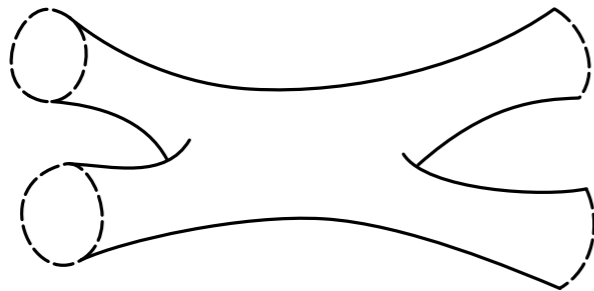
Interactions

String theory



Interactions

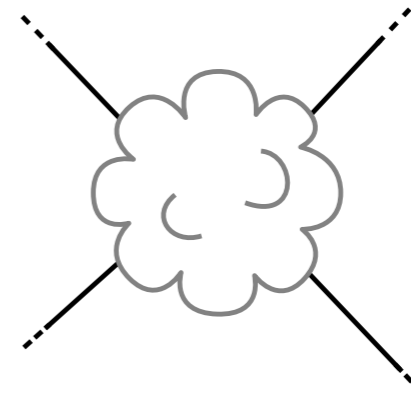
String theory



Effective field theory

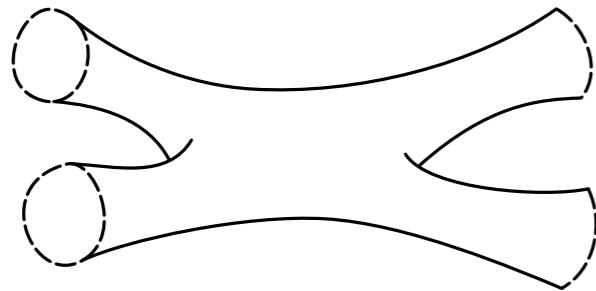


Supergravity + $\mathcal{O}(\alpha')$



Interactions

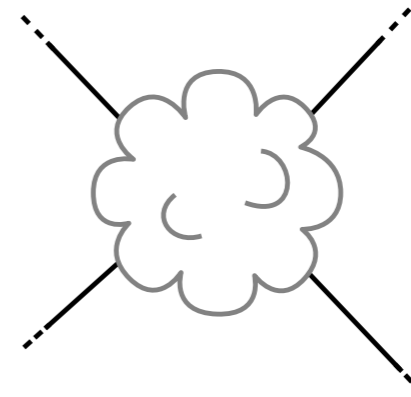
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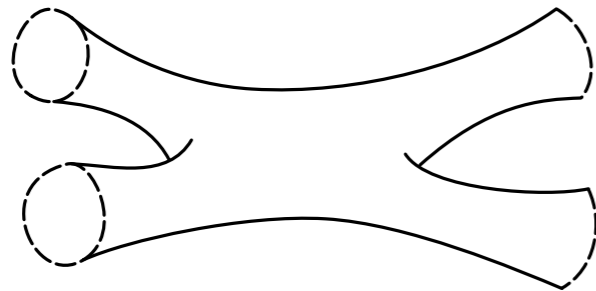
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$$s, t, u = \mathcal{O}(\alpha') \quad p \longleftrightarrow \partial \implies \alpha' \text{-expansion} = \partial \text{-expansion}$$

Interactions

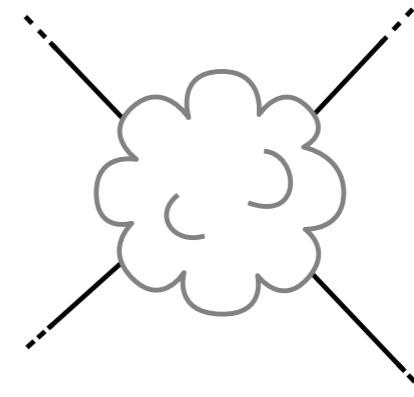
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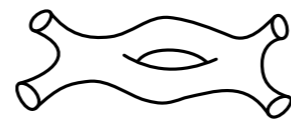
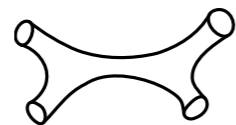
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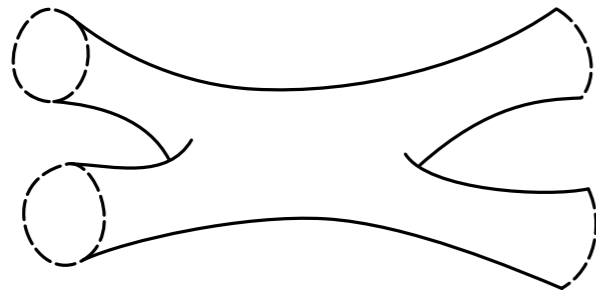


(Einstein frame)

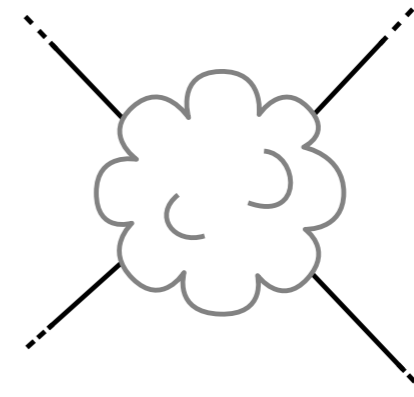
$$\mathcal{L} \propto R + (\alpha')^3 \left(2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + \dots \right) R^4 + \dots$$

Interactions

String theory



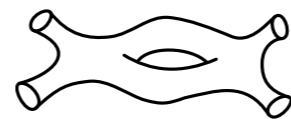
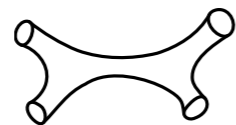
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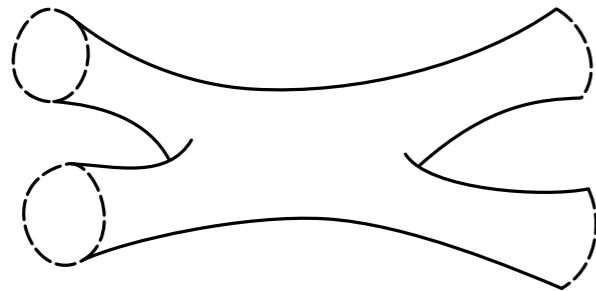
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Contraction of 4 Riemann tensors

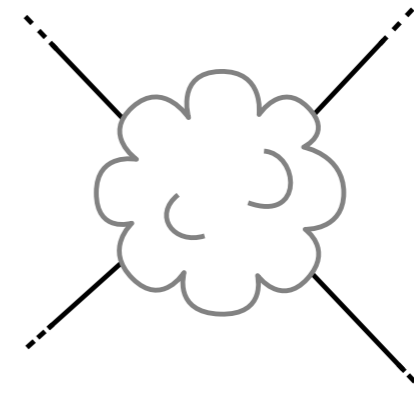


Interactions

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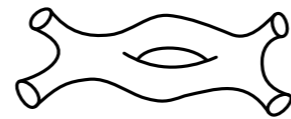
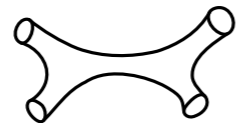
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$$\begin{aligned} \mathcal{L} \propto R + (\alpha')^3 & \left(2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + \dots \right) R^4 + \\ & (\alpha')^5 \left(\zeta(5)g_s^{-5/2} + \dots \right) D^4 R^4 + \\ & (\alpha')^6 \left(\frac{2}{3}\zeta(3)^2 g_s^{-3} + \frac{4}{3}\zeta(2)\zeta(3)g_s^{-1} + \dots \right) D^6 R^4 + \mathcal{O}((\alpha')^7) \end{aligned}$$

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$$\mathcal{L} \propto R + (\alpha')^3 \mathcal{E}_0(\tau) R^4 + (\alpha')^5 \mathcal{E}_4(\tau) D^4 R^4 + (\alpha')^6 \mathcal{E}_6(\tau) D^6 R^4 + \dots$$

$$\tau = \chi + ig_s^{-1}$$

Interactions

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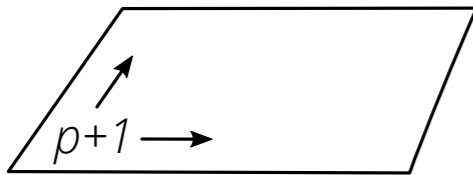
$$\tau = \chi + ig_s^{-1}$$

Ahlén, Bao, Basu, Bossard, Cederwall, Fleig, Green, Gubay, Gutperle, HG, Kazhdan, Kiritsis, Kleinschmidt, Lambert, Miller, Nilsson, Obers, Persson, Pioline, Russo, Sethi, Vanhove, Verschinin, Waldron, West, ...

Non-perturbative effects

[Green, Polchinski]

Non-perturbative effects

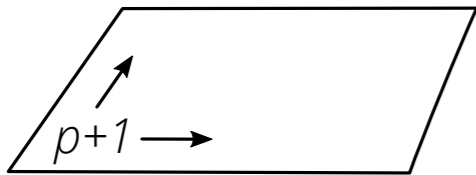


Dp -brane

p space directions

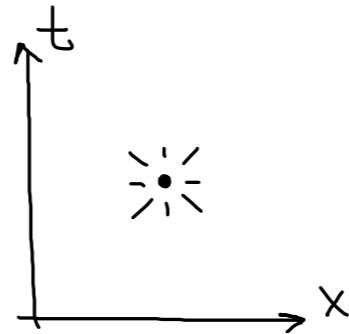
1 time direction

Non-perturbative effects



Dp -brane

p space directions
1 time direction

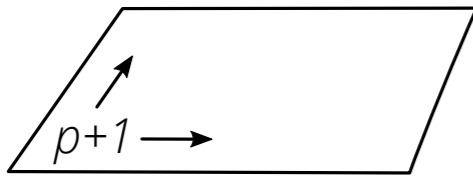


D-instanton

$p = -1$

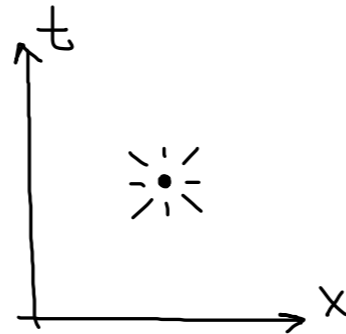
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Non-perturbative effects



Dp -brane

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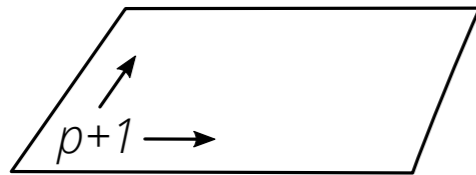


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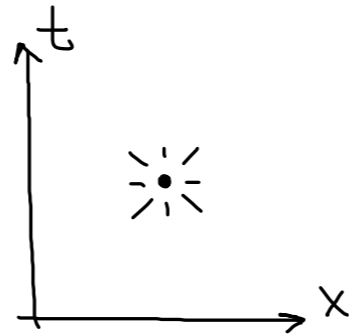


Non-perturbative effects

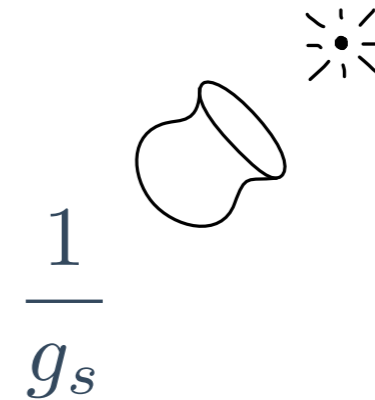


Dp -brane

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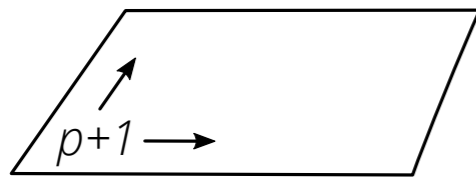


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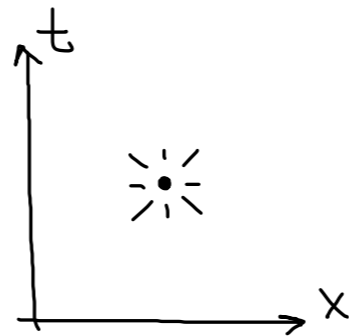
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Non-perturbative effects

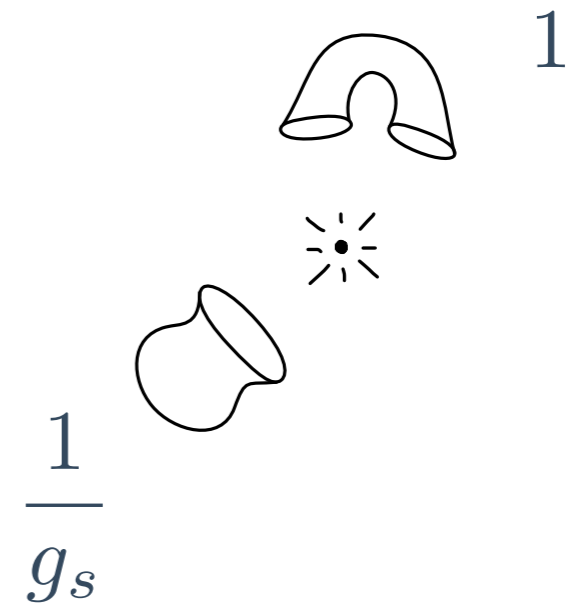


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1 time direction

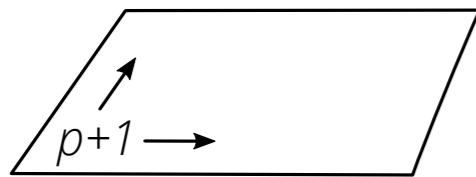


D-instanton
 $p = -1$



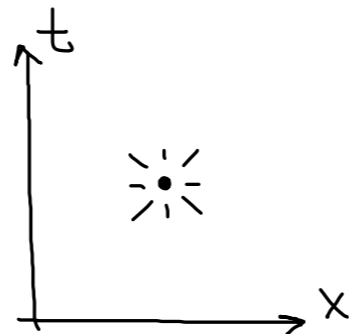
[Green, Polchinski]

Non-perturbative effects



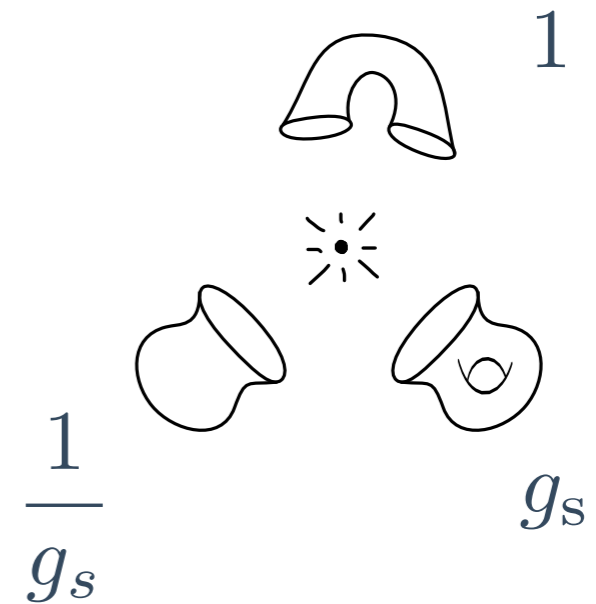
Dp -brane

p space directions
1 time direction



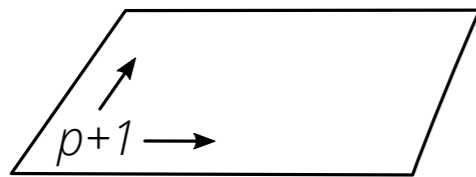
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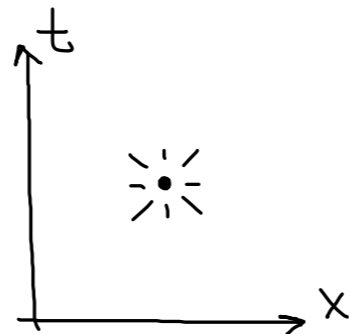
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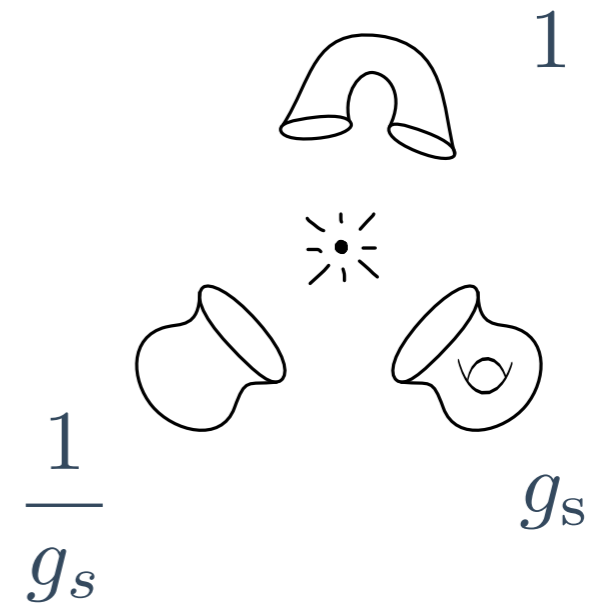
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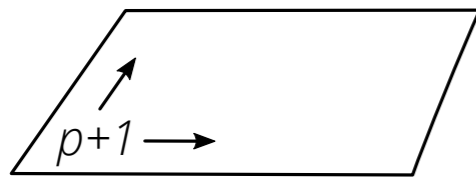
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1

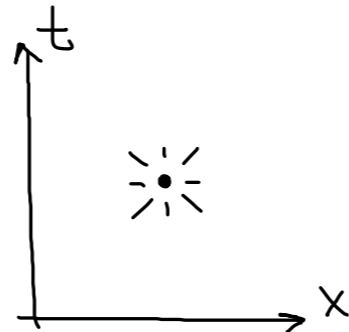
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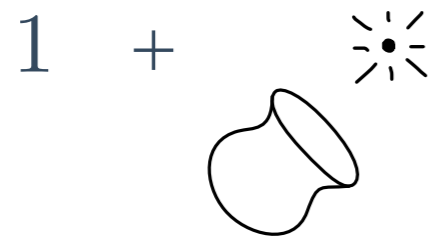
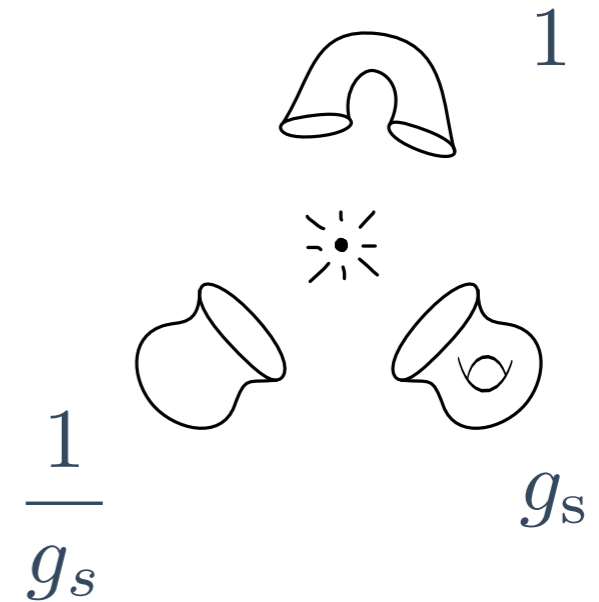


Dp -brane

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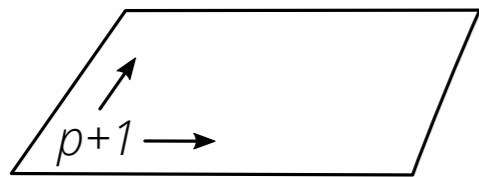


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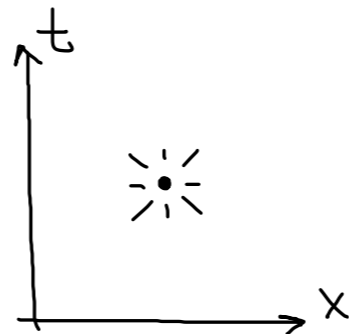
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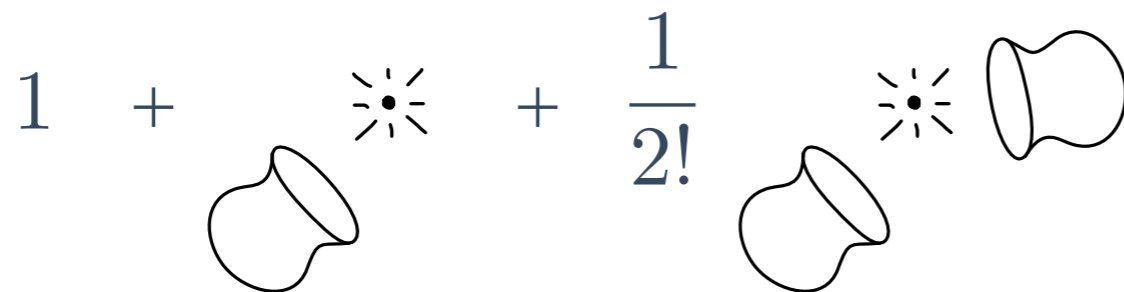
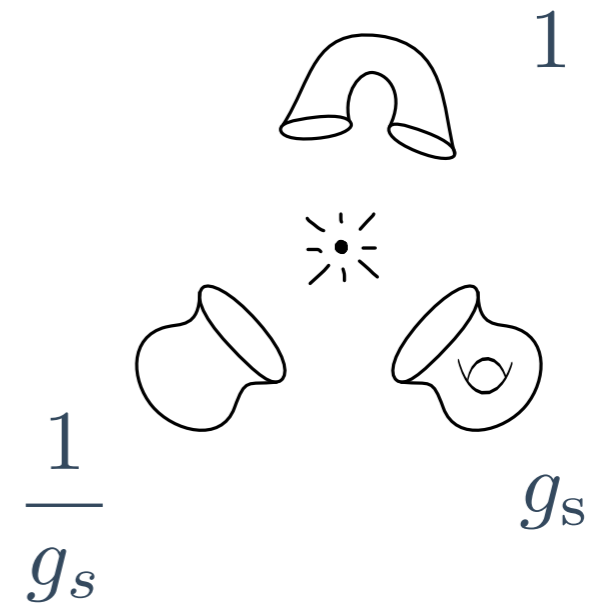


D*p*-brane

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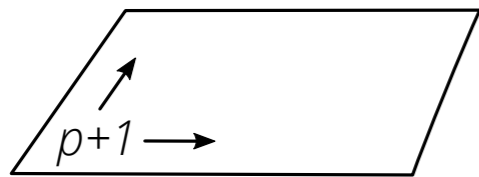


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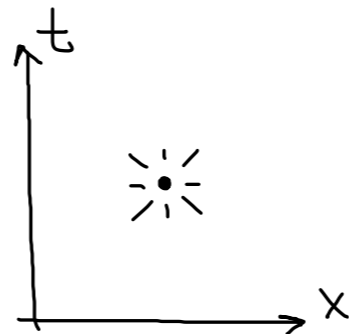
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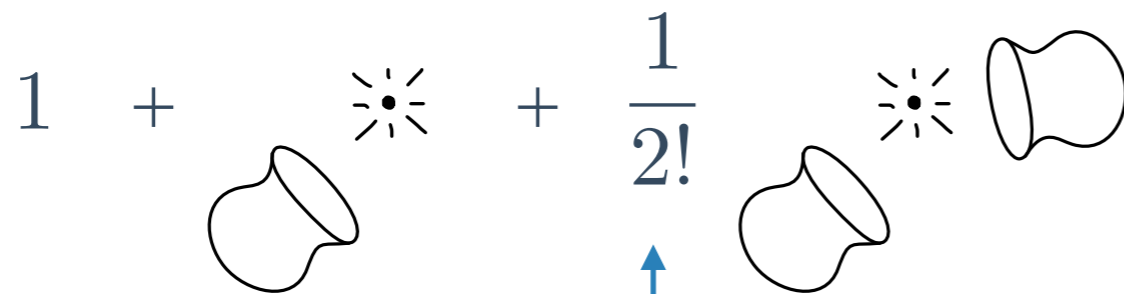
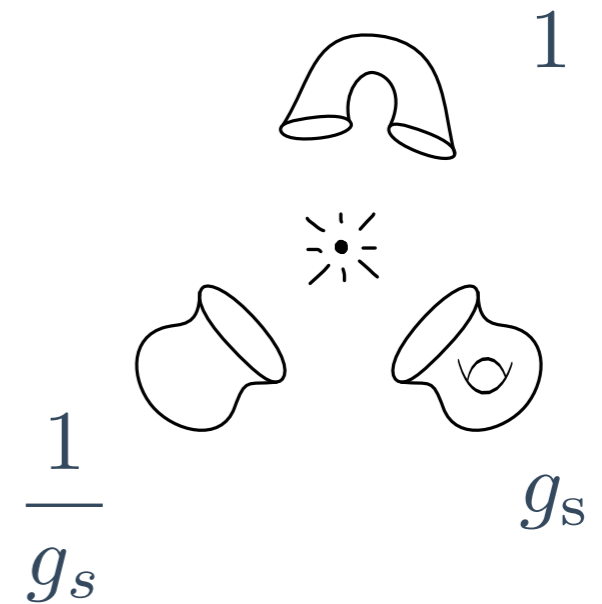


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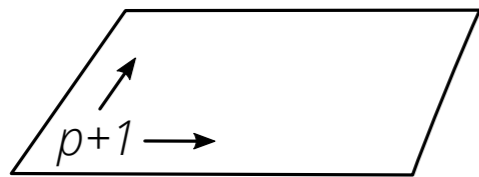
D-instanton
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Symmetry factor for identical disks

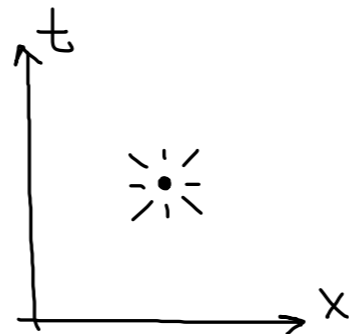
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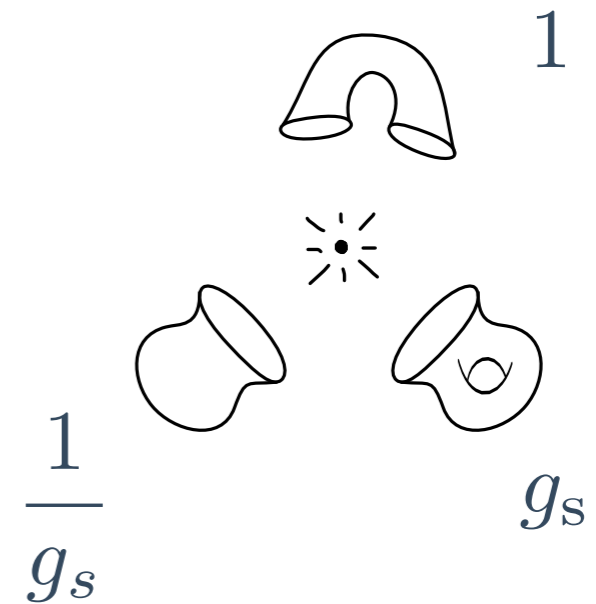


Dp -brane

p space directions
1 time direction



D-instanton
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$$1 + \text{[disk]} + \frac{1}{2!} \text{[two disks]} + \frac{1}{3!} \text{[three disks]} + \dots$$

Symmetry factor for identical disks

[Green, Polchinski]

Non-perturbative effects

$$1 + \text{[cup]} \text{[starburst]} + \frac{1}{2!} \text{[cup]} \text{[starburst]} \text{[cup]} + \frac{1}{3!} \text{[cup]} \text{[starburst]} \text{[cup]} \text{[cup]} + \dots$$

[Green-Gutperle]

Non-perturbative effects

$$1 + \text{[cup]} \text{[star]} + \frac{1}{2!} \text{[cup]} \text{[star]} \text{[cup]} + \frac{1}{3!} \text{[cup]} \text{[star]} \text{[cup]} \text{[cup]} + \dots$$
$$\exp \left(\text{[cup]} \right)$$

[Green-Gutperle]

Non-perturbative effects

$$1 + \text{[one-holed torus]} \ast + \frac{1}{2!} \text{[one-holed torus]} \ast \text{[two-holed torus]} + \frac{1}{3!} \text{[one-holed torus]} \ast \text{[two-holed torus]} \ast \text{[one-holed torus]} + \dots$$

$$\exp\left(\text{[one-holed torus]}\right) \sim \exp\left(-\frac{\text{const}}{g_s}\right)$$

Non-perturbative effects

$$1 + \text{[one-holed torus]} \cdot \text{[sunburst]} + \frac{1}{2!} \text{[one-holed torus]} \cdot \text{[sunburst]} \cdot \text{[two-holed torus]} + \frac{1}{3!} \text{[one-holed torus]} \cdot \text{[sunburst]} \cdot \text{[two-holed torus]} \cdot \text{[one-holed torus]} + \dots$$

$$\exp\left(\text{[one-holed torus]}\right) \sim \exp\left(-\frac{\text{const}}{g_s}\right) \quad \text{Non-perturbative in } g_s$$

Non-perturbative effects

$$1 + \begin{array}{c} \text{sunburst} \\ \uparrow \\ \text{torus} \end{array} + \frac{1}{2!} \begin{array}{c} \text{sunburst} \\ \uparrow \\ \text{torus} \end{array} \begin{array}{c} \text{sunburst} \\ \uparrow \\ \text{torus} \end{array} + \frac{1}{3!} \begin{array}{c} \text{sunburst} \\ \uparrow \\ \text{torus} \end{array} \begin{array}{c} \text{sunburst} \\ \uparrow \\ \text{torus} \end{array} + \dots$$

$$\exp\left(\begin{array}{c} \text{sunburst} \\ \uparrow \\ \text{torus} \end{array}\right) \sim \exp\left(-\frac{\text{const}}{g_s}\right) \quad \text{Non-perturbative in } g_s$$

$$\mathcal{E}_0(\tau) = 2\zeta(3)\tau_2^{3/2} + 4\zeta(2)\tau_2^{-1/2} + \dots + \underbrace{Ce^{2\pi i\tau}}_{\text{.....}} + \dots$$

$$\tau = \tau_1 + i\tau_2 = \chi + ig_s^{-1}$$

[Green-Gutperle]

Moduli space

$$R + (\alpha')^3 \mathcal{E}_0^{(D)}(g) R^4 + (\alpha')^5 \mathcal{E}_4^{(D)}(g) D^4 R^4 + (\alpha')^6 \mathcal{E}_6^{(D)}(g) D^6 R^4 + \dots$$

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D	$G(\mathbb{R})$	K
10	$SL(2, \mathbb{R})$	$SO(2)$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$
7	$SL(5, \mathbb{R})$	$SO(5)$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5))/\mathbb{Z}_2$
5	$E_6(\mathbb{R})$	$USp(8)/\mathbb{Z}_2$
4	$E_7(\mathbb{R})$	$SU(8)/\mathbb{Z}_2$
3	$E_8(\mathbb{R})$	$Spin(16)/\mathbb{Z}_2$

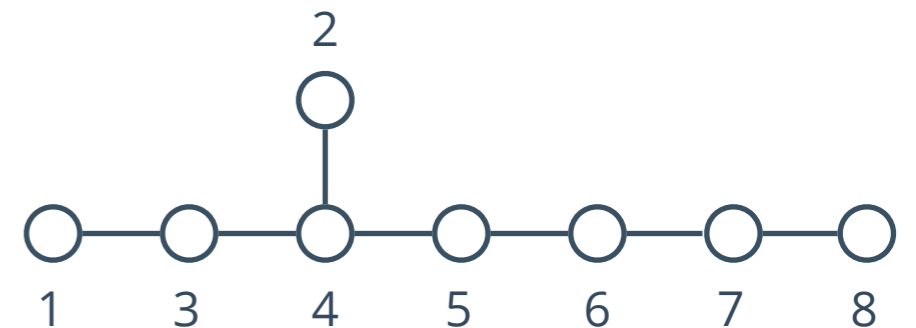
[Cremmer-Julia]

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$$\tau = \chi + ig_s^{-1} \in \mathbb{H} = \{z \in \mathbb{C} \mid \text{Im } z > 0\} \cong SL(2, \mathbb{R})/SO(2, \mathbb{R})$$

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No similar structure for lower dimensions

U-duality

$G(\mathbb{R}) \curvearrowright \mathcal{M}_{\text{classical}}$ classical symmetry

[Hull-Townsend]

U-duality

$G(\mathbb{R}) \curvearrowright \mathcal{M}_{\text{classical}}$ classical symmetry

Quantization of charges

[Hull-Townsend]

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Quantization of charges \implies classical symmetry \longrightarrow discrete symmetry

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All observables are invariant under $G(\mathbb{Z})$

[Hull-Townsend]

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$$\mathcal{E}_0^{(D)}(g), \mathcal{E}_4^{(D)}(g), \mathcal{E}_6^{(D)}(g) : G(\mathbb{Z}) \backslash G(\mathbb{R}) / K \rightarrow \mathbb{R}$$

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- (C) φ is an eigenfunction to all G -invariant differential operators

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- (C) Z-finiteness: $\dim(\text{span}\{X\varphi(g) \mid X \in \mathcal{Z}(\mathfrak{g}_{\mathbb{C}})\}) < \infty$

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- (D) Growth: for any norm $\|\cdot\|$ on $G(\mathbb{R})$ there exists a positive integer n and constant C such that $|\varphi(g)| \leq C\|g\|^n$

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Automorphic forms

An *automorphic form* is a smooth function $\varphi : G(\mathbb{R}) \rightarrow \mathbb{C}$ satisfying the following conditions

- (A) Automorphic invariance: ✓ U-duality
- (B) K-finiteness:
- (C) Z-finiteness:
- (D) Growth:

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- (B) K-finiteness: ✓ spherical
- (C) Z-finiteness:
- (D) Growth:

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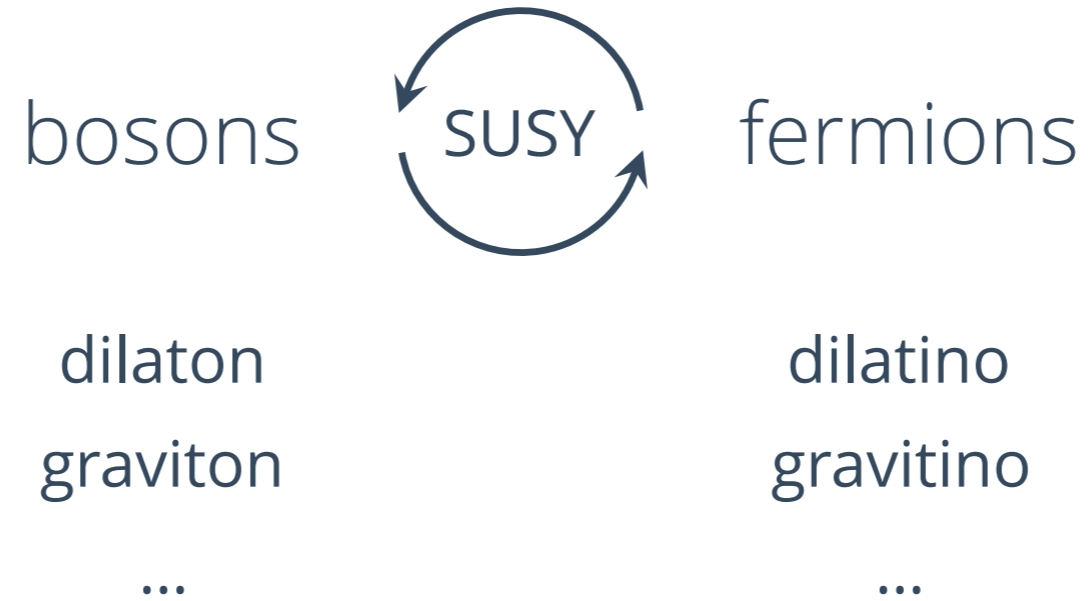
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- (D) Growth: ✓ weak coupling limit from string perturbation theory

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Supersymmetry constraints



Supersymmetry constraints



10 dimensions at order $(\alpha')^3$:

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Supersymmetry constraints



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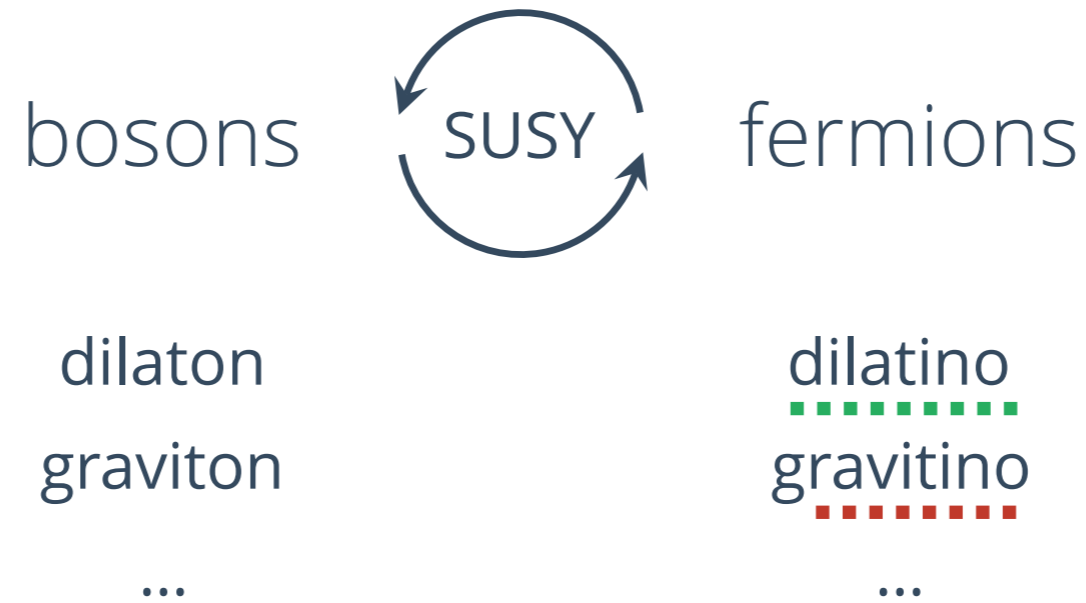
Supersymmetry constraints



10 dimensions at order $(\alpha')^3$:

$$\mathcal{L}^{(3)} = f_{12}(\tau)\lambda^{\dots 16} + f_{11}(\tau)\hat{G}\lambda^{14} + \dots + f_0(\tau)R^4 + \dots + f_{-12}(\tau)\lambda^{\dots *16}$$

Supersymmetry constraints

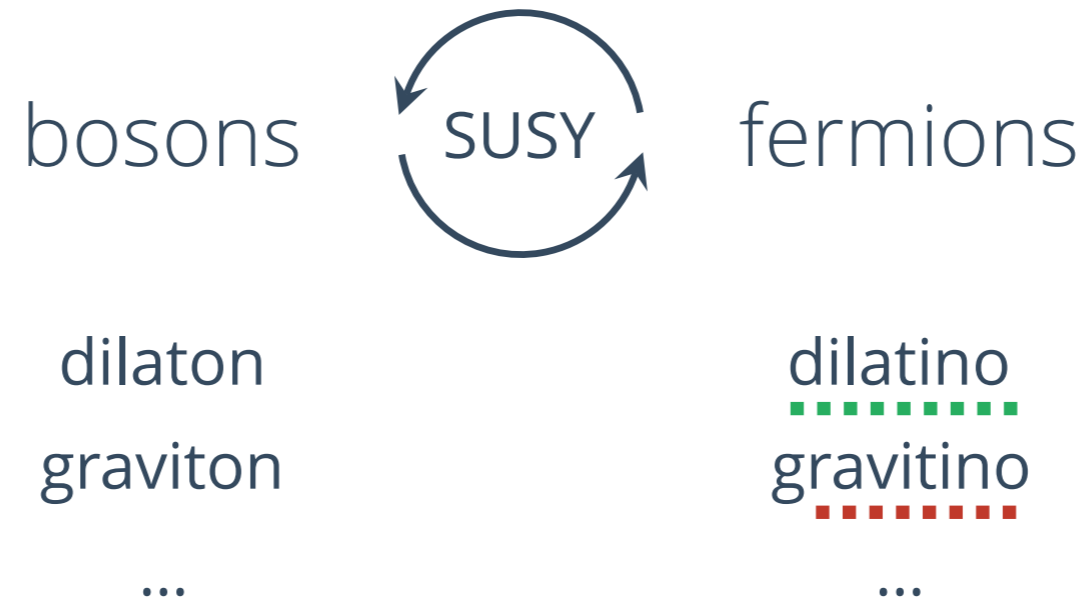


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.....

Supersymmetry constraints



10 dimensions at order $(\alpha')^3$:

$$\mathcal{L}^{(3)} = f_{12}(\tau)\lambda^{16} + f_{11}(\tau)\hat{G}\lambda^{14} + \dots + f_0(\tau)R^4 + \dots + f_{-12}(\tau)\lambda^{*16}$$

Linearized SUSY: $f_{w+1}(\tau) = i\left(\tau_2\frac{\partial}{\partial\tau} - i\frac{w}{2}\right)f_w(\tau)$

[Green-Sethi]

Supersymmetry constraints

$$\int d^D x \sqrt{-g} \mathcal{L} = S = S^{(0)}$$
$$\delta \Psi = \delta^{(0)} \Psi$$

Supersymmetry constraints

$$\int d^D x \sqrt{-g} \mathcal{L} = S = S^{(0)} + (\alpha')^3 S^{(3)} + (\alpha')^5 S^{(5)} + \dots$$
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$f_w(\tau)$ is contained in $S^{(3)}$

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$$\left\{ \begin{array}{l} \delta^{(0)} S^{(3)} + \delta^{(3)} S^{(0)} = 0 \\ [\delta_1, \delta_2] \lambda^* = \delta_{\text{local symmetries}} \lambda^* + (\text{equations of motion}) \end{array} \right.$$

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$$\implies (\Delta - \frac{3}{4}) \mathcal{E}_0(\tau) = 0$$

[Green-Sethi]

Supersymmetry constraints

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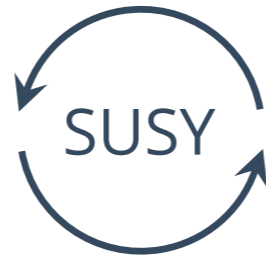
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$$\implies \left(\Delta - \frac{3}{4} \right) \mathcal{E}_0(\tau) = 0 \quad \Delta = 4\tau_2^2 \frac{\partial}{\partial \tau} \frac{\partial}{\partial \bar{\tau}} \quad \text{Laplacian on Poincaré UHP}$$

[Green-Sethi]

Supersymmetry constraints

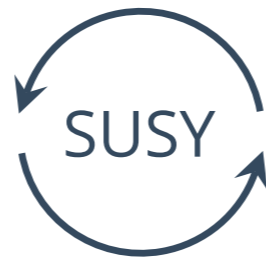


10 dimensions: $\Delta = 4\tau_2^2 \frac{\partial}{\partial \tau} \frac{\partial}{\partial \bar{\tau}}$ Laplacian on
Poincaré UHP

$$\left(\Delta - \frac{3}{4}\right) \mathcal{E}_0(\tau) = 0$$

[Green-Sethi]

Supersymmetry constraints



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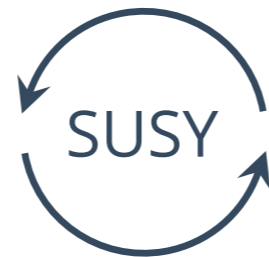
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Supersymmetry constraints



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$$(\Delta - 12)\mathcal{E}_6(\tau) = -(\mathcal{E}_0(\tau))^2$$

[Green-Vanhove]

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Not an automorphic form in a strict sense

Supersymmetry constraints



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Not an automorphic form in a strict sense

Similarly for lower dimensions

Eisenstein series

$$E(s; \tau) =$$

$$s \in \mathbb{C}$$

Eisenstein series

$$E(s; \tau) = \sum_{\substack{c, d \in \mathbb{Z} \\ (c, d) \neq (0, 0)}} \frac{\text{Im}(\tau)^s}{|c\tau + d|^{2s}}$$

$s \in \mathbb{C}$

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$$E(s; \tau) = \frac{1}{2\zeta(2s)} \sum_{\substack{c, d \in \mathbb{Z} \\ (c, d) \neq (0, 0)}} \frac{\text{Im}(\tau)^s}{|c\tau + d|^{2s}}$$

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$$E(s; \tau) = \frac{1}{2\zeta(2s)} \sum_{\substack{c, d \in \mathbb{Z} \\ (c, d) \neq (0, 0)}} \frac{\text{Im}(\tau)^s}{|c\tau + d|^{2s}} = \sum_{\gamma \in B(\mathbb{Z}) \setminus SL(2, \mathbb{Z})} \frac{\text{Im}(\tau)^s}{|c\tau + d|^{2s}} \quad s \in \mathbb{C}$$

$$B(\mathbb{Z}) = \left\{ \begin{pmatrix} * & * \\ 0 & * \end{pmatrix} \in SL(2, \mathbb{Z}) \right\} \quad \text{Borel subgroup}$$

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Multiplicative character
trivially extended to $G(\mathbb{R})$

$$\tau \mapsto \text{Im}(\tau)^s$$

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of the second kind

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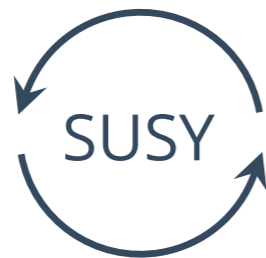
Eisenstein series

$$(\Delta - s(s - 1))E(s; \tau) = 0 \quad E(s; \tau) \sim \tau_2^s \quad g_s = \tau_2^{-1} \rightarrow 0$$

[Green-Gutperle, Pioline, Green-Russo-Vanhove]

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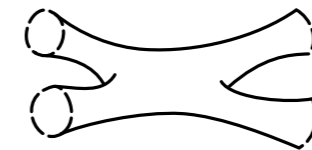
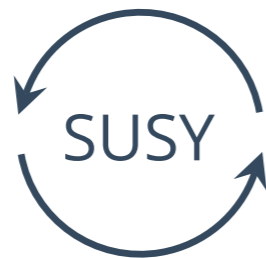
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(Einstein frame)

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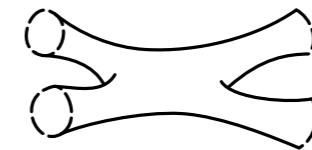
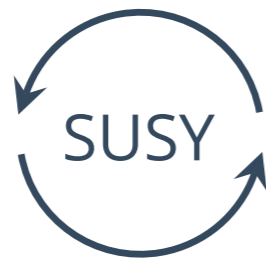
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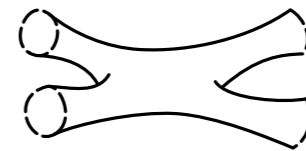
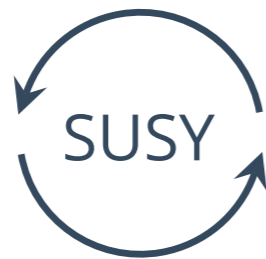
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[Green-Gutperle, Pioline, Green-Russo-Vanhove]

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$\mathcal{E}_6(\tau)$ as a sum over images $\sum_{B(\mathbb{Z}) \backslash G(\mathbb{Z})}$ but not of a character χ

[Green-Miller-Vanhove]

Extracting physical information

Expand Bessel function in g_s

$$\tau = \chi + ig_s^{-1}$$

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Perturbative
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[Green-Gutperle]

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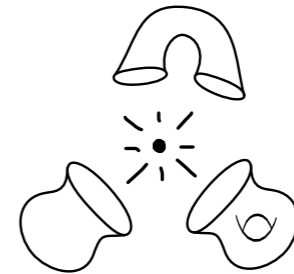
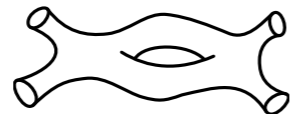
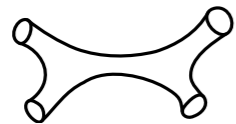
Perturbative
(zero-mode)

Non-perturbative
(remaining modes)

Extracting physical information

Expand Bessel function in g_s

$$\tau = \chi + ig_s^{-1}$$



$$\mathcal{E}_0(\tau) = 2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-2\pi|m|g_s^{-1} + 2\pi im\chi} \left[1 + \mathcal{O}(g_s) \right]$$

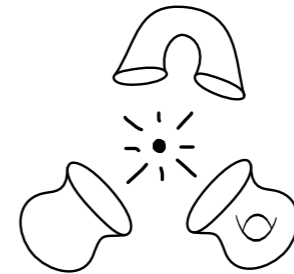
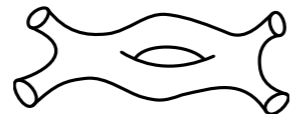
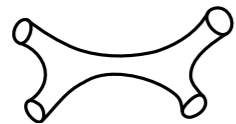
Perturbative
(zero-mode)

Non-perturbative
(remaining modes)

Extracting physical information

Expand Bessel function in g_s

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Instanton action

$$\mathcal{E}_0(\tau) = 2\zeta(3)g_s^{-3/2} + 4\zeta(2)g_s^{1/2} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-2\pi|m|g_s^{-1} + 2\pi im\chi} \left[1 + \mathcal{O}(g_s) \right]$$

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Instanton action



Perturbative
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Non-perturbative
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$$\sigma_s(m) = \sum_{d|m} d^s$$

Sums over the number of ways the charge m can be factorised into two integers

Extracting physical information

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$$\sigma_s(m) = \sum_{d|m} d^s$$

Sums over the number of ways the charge m can be factorised into two integers

wrapping number and charge
of a T-dual D-particle

[Green-Gutperle]

Lower dimensions

Lower dimensions

D	$G(\mathbb{R})$	K	$G(\mathbb{Z})$
10	$SL(2, \mathbb{R})$	$SO(2)$	$SL(2, \mathbb{Z})$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$	$SL(2, \mathbb{Z}) \times \mathbb{Z}_2$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$	$SL(3, \mathbb{Z}) \times SL(2, \mathbb{Z})$
7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5; \mathbb{R})$	$(Spin(5) \times Spin(5)) / \mathbb{Z}_2$	$Spin(5, 5; \mathbb{Z})$
5	$E_6(\mathbb{R})$	$USp(8) / \mathbb{Z}_2$	$E_6(\mathbb{Z})$
4	$E_7(\mathbb{R})$	$SU(8) / \mathbb{Z}_2$	$E_7(\mathbb{Z})$
3	$E_8(\mathbb{R})$	$Spin(16) / \mathbb{Z}_2$	$E_8(\mathbb{Z})$

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7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
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$$E(\chi; g) = \sum_{\gamma \in B(\mathbb{Z}) \setminus G(\mathbb{Z})} \chi(\gamma g)$$

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

$\langle \Sigma \rangle$ generated root system

Parabolic subgroups

Fourier expand
in different directions



Unipotent subgroup U



Choice of parabolic subgroup P

Σ choice of simple roots

$\langle \Sigma \rangle$ generated root system

$$\mathfrak{g}_\alpha = \{g \in \mathfrak{g} \mid [h, g] = \alpha(h)g \quad \forall h \in \mathfrak{h}\}$$

Cartan subalgebra 

Parabolic subgroups

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Unipotent subgroup U



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Cartan subalgebra \nearrow

$$\mathfrak{p} = \mathfrak{l} \oplus \mathfrak{u}$$

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$$\mathfrak{p} = \mathfrak{l} \oplus \mathfrak{u}$$

$$\mathfrak{l} = \mathfrak{h} \oplus \bigoplus_{\alpha \in \langle \Sigma \rangle} \mathfrak{g}_\alpha$$

$$\mathfrak{u} = \bigoplus_{\alpha \in \Delta(\mathfrak{u})} \mathfrak{g}_\alpha$$

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$$\Delta(\mathfrak{u}) = \Delta_+ \setminus (\Delta_+ \cap \langle \Sigma \rangle)$$

\nwarrow Positive roots

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Fourier expand
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$$G = SL(4)$$



$$\Sigma = \{\alpha_1\}$$

Parabolic subgroups

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$$L = \left\{ \begin{pmatrix} * & * & & \\ * & * & & \\ & & * & \\ & & & * \end{pmatrix} \right\}$$

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Parabolic subgroups

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$$P = LU = \left\{ \begin{pmatrix} * & * & * & * \\ * & * & * & * \\ & * & * & * \\ & & * & * \end{pmatrix} \right\} \quad L = \left\{ \begin{pmatrix} * & * & & \\ * & * & & \\ & * & & \\ & & * & \end{pmatrix} \right\} \quad U = \left\{ \begin{pmatrix} 1 & * & * & \\ & 1 & * & * \\ & & 1 & * \\ & & & 1 \end{pmatrix} \right\}$$

Parabolic subgroups



$$L = \left\{ \begin{pmatrix} * & * & & & \\ * & * & & & \\ & & * & & \\ & & & * & \\ & & & & * \end{pmatrix} \right\} \quad U = \left\{ \begin{pmatrix} 1 & & * & * & \\ & 1 & * & * & \\ & & 1 & * & \\ & & & 1 & * \\ & & & & 1 \end{pmatrix} \right\}$$

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Parabolic subgroups



Parabolic subgroups



Maximal parabolic

Parabolic subgroups



Minimal parabolic
Borel



Maximal parabolic

Parabolic subgroups



Minimal parabolic
Borel

$$B = NA$$

$$N = \left\{ \begin{pmatrix} \boxed{1} & * & * & * \\ & \boxed{1} & * & * \\ & & \boxed{1} & * \\ & & & \boxed{1} \end{pmatrix} \right\}$$



Maximal parabolic

$$P = LU$$

$$U = \left\{ \begin{pmatrix} \boxed{1} & & & * \\ & \boxed{1} & & * \\ & & \boxed{1} & * \\ & & & \boxed{1} \end{pmatrix} \right\}$$

Fourier expansion

Fourier expansion

Let $\psi : U(\mathbb{Z}) \backslash U(\mathbb{R}) \rightarrow U(1)$ be a multiplicative character $\psi(u_1 u_2) = \psi(u_1) \psi(u_2)$

Parametrised by $m_\alpha \in \mathbb{Z}$ called charges

Fourier expansion

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$$u = \prod_{\alpha \in \Delta^{(1)}(\mathfrak{u})} \exp(u_\alpha E_\alpha) \mapsto \exp\left(2\pi i \sum_{\alpha \in \Delta^{(1)}(\mathfrak{u})} m_\alpha u_\alpha\right)$$

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$$\psi_U \left(\begin{pmatrix} 1 & & & \\ & 1 & y_1 & \\ & & 1 & y_2 \\ & & & 1 & y_3 \\ & & & & 1 \end{pmatrix} \right) = e^{2\pi i(m_1 y_1 + m_2 y_2 + m_3 y_3)}$$

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Fourier expansion

$$\psi(u_1 u_2) = \psi(u_1) \psi(u_2)$$

Let $\psi : U(\mathbb{Z}) \backslash U(\mathbb{R}) \rightarrow U(1)$ be a multiplicative character

Parametrised by $m_\alpha \in \mathbb{Z}$ called charges



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$$F_U(\chi, \psi; g) = \int_{U(\mathbb{Z}) \backslash U(\mathbb{R})} E(\chi, ug) \overline{\psi(u)} du$$

Fourier expansion

Fourier expansion

$$E(\chi; g) = \sum_{\psi} F_U(\chi, \psi; g)$$

Fourier expansion

$$E(\chi; g) = F_U(\chi, 1; g) + \sum_{\psi \neq 1} F_U(\chi, \psi; g)$$

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Fourier expansion

$$E(\chi; g) = F_U(\chi, 1; g) + \sum_{\psi^{(1)} \neq 1} F_{U^{(1)}}(\chi, \psi^{(1)}; g) + \sum_{\psi^{(2)} \neq 1} F_{U^{(2)}}(\chi, \psi^{(2)}; g) + \dots$$

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$$F_U(\chi, \psi; ug) = \psi(u)F_U(\chi, \psi; g) \quad \psi(u_1 u_2) = \psi(u_1)\psi(u_2)$$

$$U^{(1)} = U \quad U^{(n+1)} = [U^{(n)}, U^{(n)}]$$

Fourier expansion

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$$N : \quad \psi^{(1)} \left(\begin{pmatrix} 1 & x_1 & * & * \\ & 1 & x_2 & * \\ & & 1 & x_3 \\ & & & 1 \end{pmatrix} \right)$$

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$$F_U(\chi, \psi; ug) = \psi(u) F_U(\chi, \psi; g) \quad \psi(u_1 u_2) = \psi(u_1) \psi(u_2)$$

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Terminology

$P = B \longrightarrow U = N$ Fourier coefficient is a Whittaker coefficient

$$N = \left\{ \begin{pmatrix} 1 & * & * & * \\ & 1 & * & * \\ & & 1 & * \\ & & & 1 \end{pmatrix} \right\}$$

Terminology

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F_U

W_N

Terminology

$P = B \longrightarrow U = N$ Fourier coefficient is a Whittaker coefficient

$$N = \left\{ \begin{pmatrix} 1 & * & * & * \\ & 1 & * & * \\ & & 1 & * \\ & & & 1 \end{pmatrix} \right\} \quad F_U \quad W_N$$

Characters and coefficients with all $m_\alpha \neq 0$ are called **generic**
otherwise they are called **degenerate**

Fourier expansion

Choice of unipotent subgroup U \longleftrightarrow Study different perturbative and non-perturbative effects

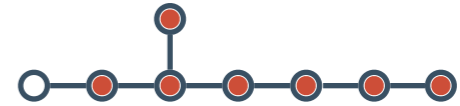
[Green-Miller-Vanhove]

Fourier expansion

Choice of unipotent subgroup U \longleftrightarrow Study different perturbative and non-perturbative effects

- String perturbation limit
D-instantons | NS5-instantons

$$g_s \rightarrow 0$$



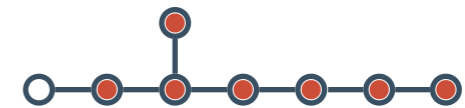
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Fourier expansion

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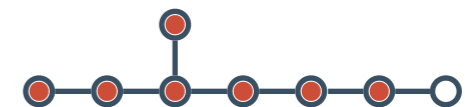
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$$g_s \rightarrow 0$$



- Decompactification limit
Higher dimensional black holes | BPS states

Large radius for
compactified circle



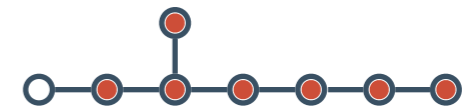
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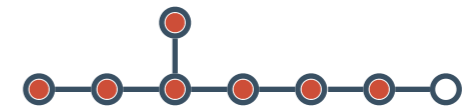
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Large radius for
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- M-theory limit
M2, M5-instantons

Large M-theory torus



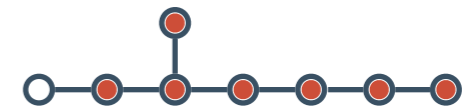
[Green-Miller-Vanhove]

Fourier expansion

Choice of unipotent subgroup U \longleftrightarrow Study different perturbative and non-perturbative effects

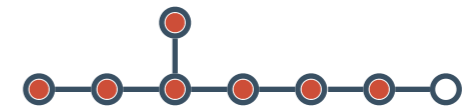
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$$g_s \rightarrow 0$$



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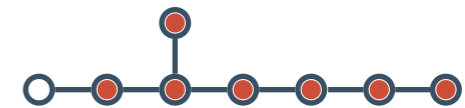
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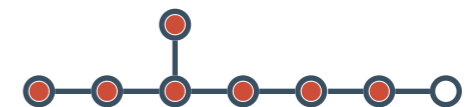
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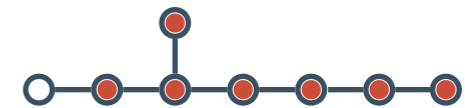
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[Green-Miller-Vanhove]

Maximal parabolic subgroups

Difficult to compute!

Recent result in [Bossard-Pioline]

Fourier expansion

Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients

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Would allow us to compute non-perturbative effects that
capture information about instantons and black holes

Adelic framework

*An **efficient**, but abstract, way to approach the subject of automorphic forms is by the introduction of **adeles**, rather **ungainly objects** that nevertheless, once familiar, **spare** much unnecessary thought and **many useless calculations**.*

— Robert P. Langlands*

*Representation theory - its rise and its role in number theory, Proceedings of the Gibbs symposium (1989)

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Product of Cauchy completions of \mathbb{Q}

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Automorphic form : $G(\mathbb{Z}) \backslash G(\mathbb{R}) / K \rightarrow \mathbb{C}$

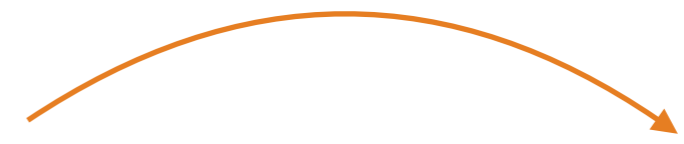
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[FGKP15 §4.2.2]



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Charges / Fourier modes
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Computing adelic Fourier coefficients

[FGKP15 §9-10]

W_N Whittaker coefficients

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Factorisation

Computing adelic Fourier coefficients

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[GKP14] + ...

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~~Factorisation~~

Write in terms of Whittaker coefficients

Simplify drastically for certain Eisenstein series, or χ

Example of simplifications

$$G = SL(3)$$

$$E(\chi; g)$$

$$\chi \longleftrightarrow (s_1, s_2) \in \mathbb{C}^2$$

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$$\psi_{m_1, m_2} \left(\begin{pmatrix} 1 & x_1 & * \\ & 1 & x_2 \\ & & 1 \end{pmatrix} \right) = e^{2\pi i(m_1 x_1 + m_2 x_2)}$$

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Vanishes for certain (s_1, s_2)

[FGKP15 §10.6]

Example of simplifications

Certain (s_1, s_2)

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To explain this, we need to study
small automorphic representations

[FGKP15 §10.6]

Automorphic representations

$G(\mathbb{A}) \curvearrowright$ Space of automorphic forms*

* With some subtleties described in [FGKP15 §6]

[Bump, Goldfeld-Hundley]

Automorphic representations

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Automorphic representation π = an irreducible component of the above space under this action

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What is a small automorphic representation?

* With some subtleties described in [FGKP15 §6]

[Bump, Goldfeld-Hundley]

Wavefront set

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg,
Gomez-Gourevitch-Sahi, Jiang-Liu-Savin, Joseph, Miller-Sahi]

Wavefront set

(Fourier modes)

The (global) wavefront set contains all the characters ψ which can give rise to non-vanishing Fourier coefficients in that representation

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Small automorphic representations have few non-vanishing Fourier coefficients

[Mœglin–Waldspurger, Matumoto, Ginzburg-Rallis-Soudry, Ginzburg, Gomez-Gourevitch-Sahi, Jiang-Liu-Savin, Joseph, Miller-Sahi]

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The wavefront set is described by nilpotent orbits

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Closure with respect to partial ordering

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[Collingwood-McGovern]

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 decreasing order
 $(p_1, p_2, \dots) \leq (q_1, q_2, \dots)$ partial ordering

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[Collingwood-McGovern]

For $SL(n)$, orbits can be identified with partitions of n

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Illustrated by a Hasse diagram

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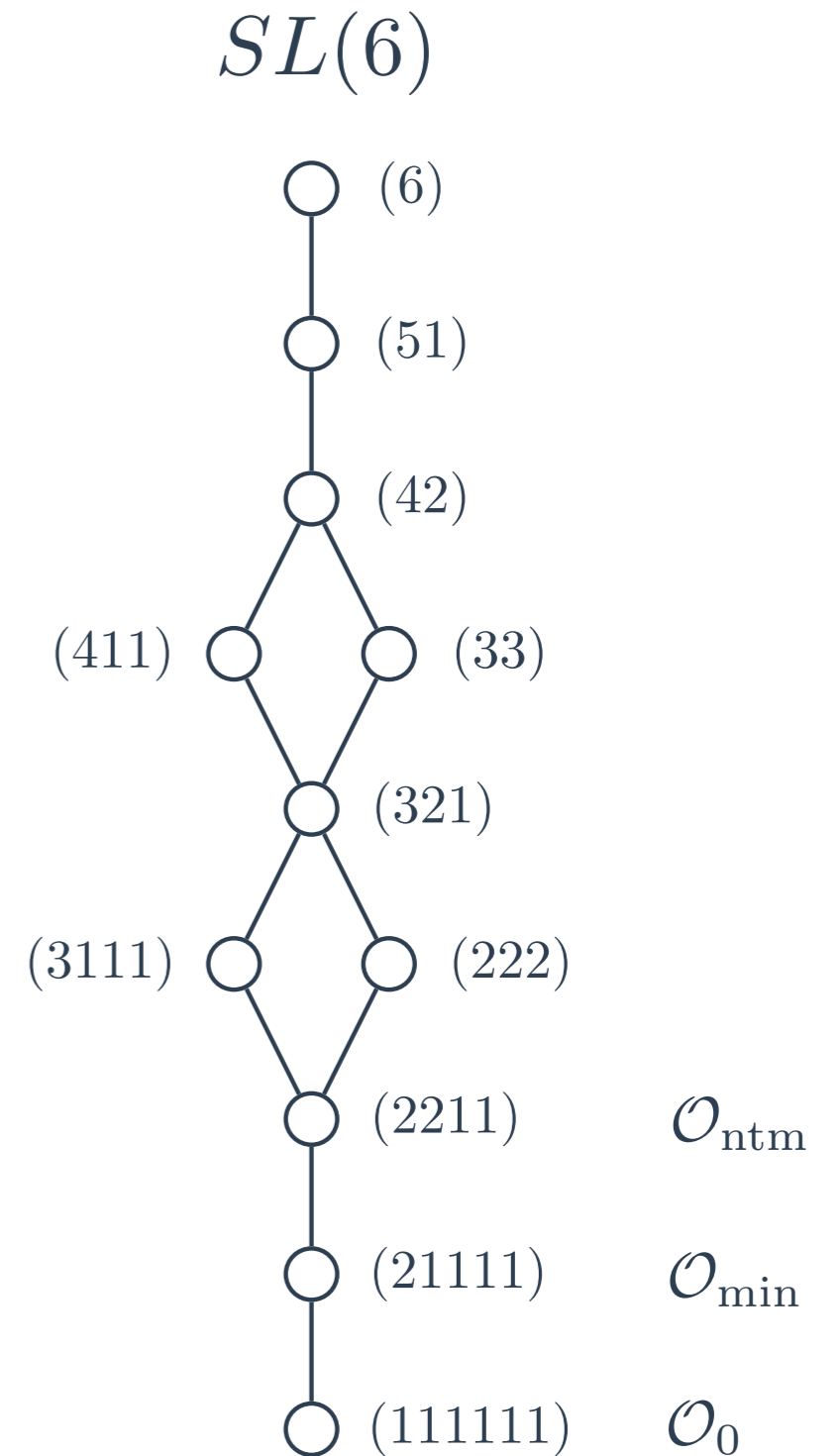
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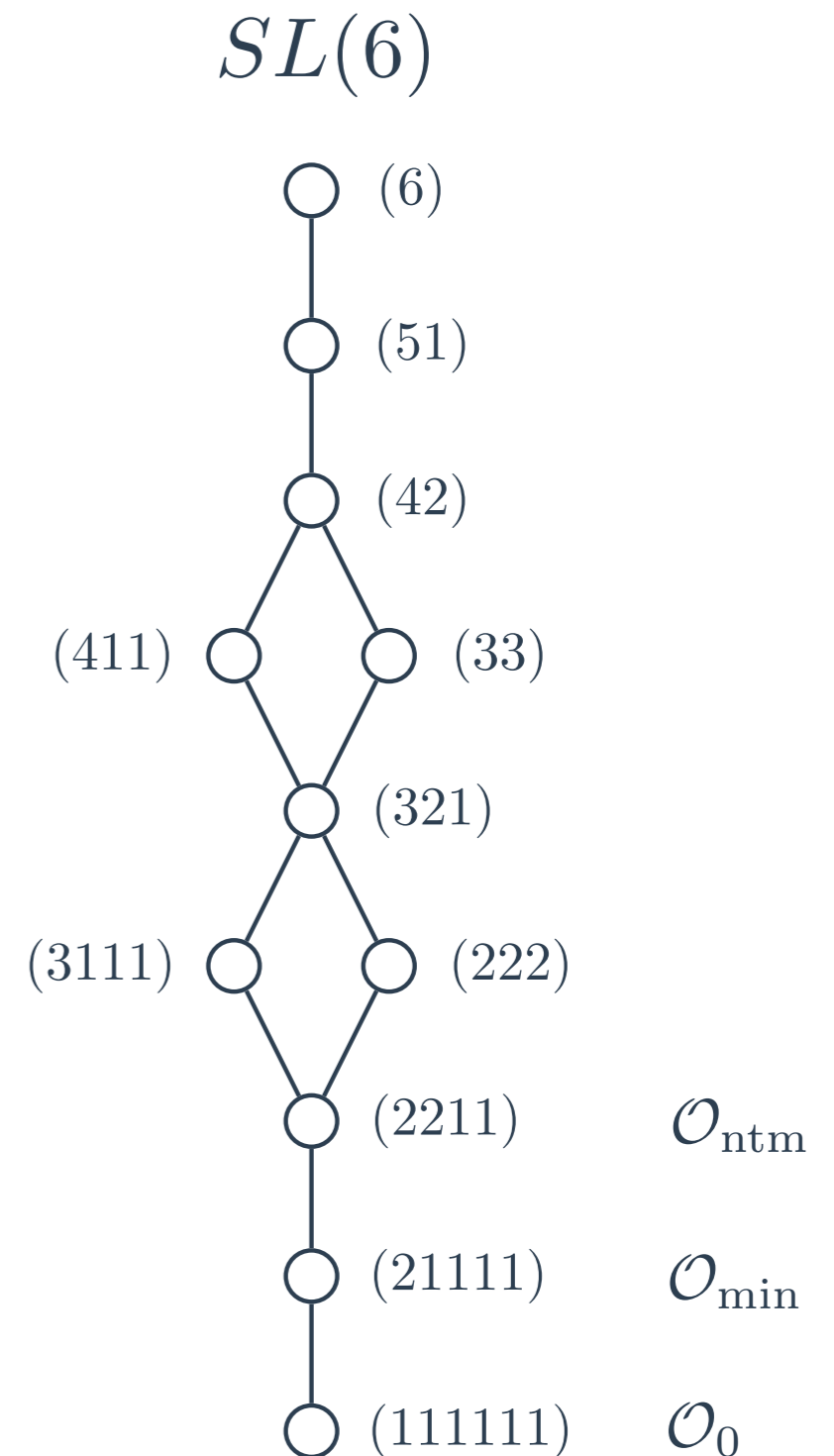
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Illustrated by a Hasse diagram

Closure: $\overline{\mathcal{O}} = \bigcup_{\mathcal{O}' \leq \mathcal{O}} \mathcal{O}'$



Automorphic representations

Small representations

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Small representations

$$\mathrm{WF}(\pi_{\min}) = \overline{\mathcal{O}_{\min}} = \mathcal{O}_{\min} \cup \mathcal{O}_0$$

$$\mathrm{WF}(\pi_{\mathrm{ntm}}) = \overline{\mathcal{O}_{\mathrm{ntm}}} = \mathcal{O}_{\mathrm{ntm}} \cup \mathcal{O}_{\min} \cup \mathcal{O}_0$$

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$$\mathcal{E}_0^{(D)} \in \pi_{\min}$$

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[Green-Miller-Vanhove,
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χ_{\min} such that $E(\chi_{\min}, g) \in \pi_{\min}$

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[Green-Miller-Vanhove,
Pioline, Bossard-Verschinin]

Certain $(s_1, s_2) \longleftrightarrow \chi_{\min}$ such that $E(\chi_{\min}, g) \in \pi_{\min}$

Automorphic representations

Small representations

$$\mathrm{WF}(\pi_{\min}) = \overline{\mathcal{O}_{\min}} = \mathcal{O}_{\min} \cup \mathcal{O}_0$$

$$\mathrm{WF}(\pi_{\mathrm{ntm}}) = \overline{\mathcal{O}_{\mathrm{ntm}}} = \mathcal{O}_{\mathrm{ntm}} \cup \mathcal{O}_{\min} \cup \mathcal{O}_0$$

$$\mathcal{E}_0^{(D)} \in \pi_{\min}$$

$$\mathcal{E}_4^{(D)} \in \pi_{\mathrm{ntm}}$$

[Green-Miller-Vanhove,
Pioline, Bossard-Verschinin]

Certain $(s_1, s_2) \longleftrightarrow \chi_{\min}$ such that $E(\chi_{\min}, g) \in \pi_{\min}$

$$\int K K \longrightarrow 0$$

$$\sum K \longrightarrow K$$

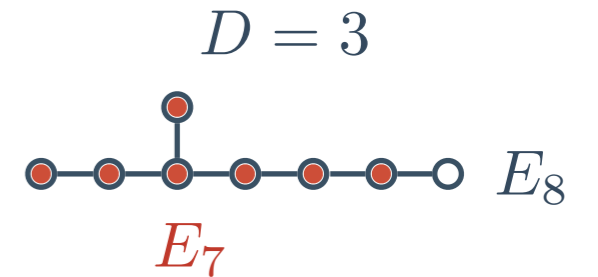
Automorphic representations

[Ferrara-Günaydin, Ferrara-Maldacena, Green-Miller-Vanhove]

Automorphic representations

- Decompactification limit
Higher dimensional black holes | BPS states

Large radius for
compactified circle

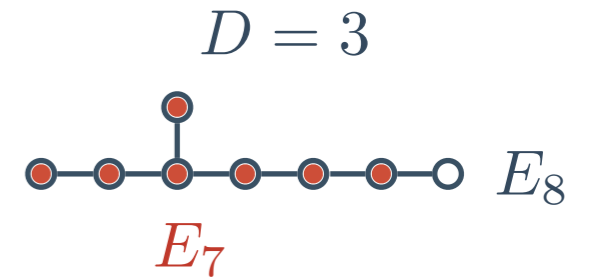


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π_{\min}

π_{ntm}

π_{3A_1}

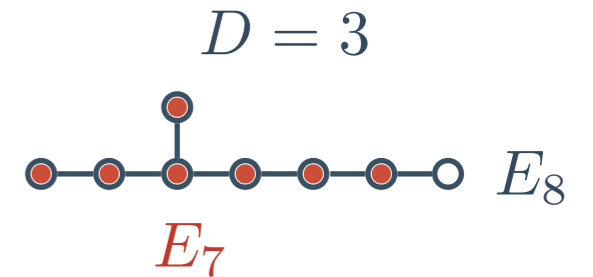
π_{A_2}

[Ferrara-Günaydin, Ferrara-Maldacena, Green-Miller-Vanhove]

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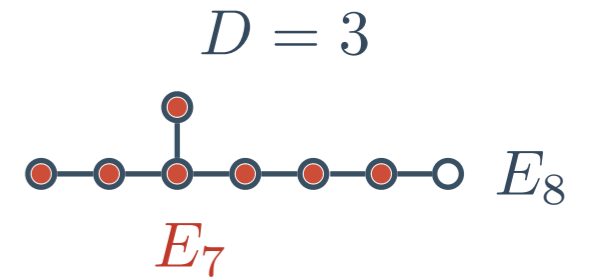
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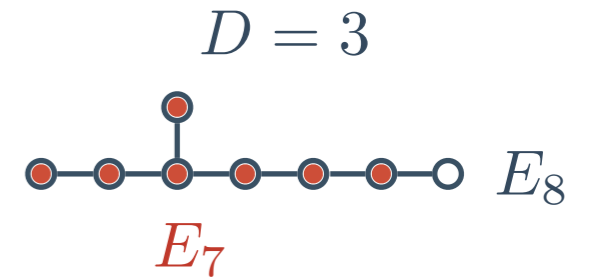
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$\dim\{\psi_U \in \text{WF}\}$	28	45	55	56	[Miller-Sahi]

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$\dim\{\psi_U \in \text{WF}\}$	28	45	55	56	[Miller-Sahi]
$D = 4$ BPS-orbits $E_7 \curvearrowright \{\psi_U\}$	$\frac{1}{2}$ BPS	$\frac{1}{4}$ BPS	$\frac{1}{8}$ BPS	$\frac{1}{8}$ BPS ⁺	

[Ferrara-Günaydin, Ferrara-Maldacena, Green-Miller-Vanhove]

Goal: find expressions for Fourier coefficients
in terms of (known) Whittaker coefficients
using vanishing properties of the given π

Previous results

[Miller-Sahi]

Previous results

Theorem

For $G = E_6, E_7$, an automorphic form $\varphi \in \pi_{\min}$ is completely determined by maximally degenerate Whittaker coefficients

W_N with only one $m_\alpha \neq 0$

[Miller-Sahi]

Main results

$SL(3), SL(4)$

[GKP14]

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For $G = SL(3), SL(4)$, an automorphic form $\varphi \in \pi_{\min}$ is completely determined by maximally degenerate Whittaker coefficients.

[GKP14]

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[GKP14]

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$$\varphi = \sum_{\mathcal{O}} \varphi_{\mathcal{O}} \quad \text{where } \varphi_{\mathcal{O}} \text{ vanishes unless } \mathcal{O} \subseteq \text{WF}(\pi)$$

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[GKP14]

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Corollary

$$\varphi \in \pi_{\min}$$

single root

$$\varphi \in \pi_{\text{ntm}}$$

at most two commuting roots

[GKP14]

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Fourier coefficients on maximal parabolic subgroups in the minimal representation



π_{\min}

[GKP14]

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$SL(3), SL(4)$

Fourier coefficients on maximal parabolic subgroups in the minimal representation



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Theorem

$$F_U(\chi_{\min}, \psi; g) = W_N(\chi_{\min}, \psi'; lg) \quad \text{with } l \in L(\mathbb{Q}) \text{ depending on } \psi$$

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Maximal parabolic
Fourier coefficient



Maximally degenerate

[GKP14]

Other groups

$$SL(n)$$

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↑
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↑
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[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

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↑
Maximal parabolic
Fourier coefficient

↑
Maximally degenerate

and similar statement for next-to-minimal representation

[Work in progress with Ahlén, Liu, Kleinschmidt, Persson]

Other groups

Conjecture

A similar relations holds for all simple Lie groups

$$F_U(\chi_{\min}, \psi; g) = W_N(\chi_{\min}, \psi'; lg) \quad \text{with } l \in L(\mathbb{Q}) \text{ depending on } \psi$$

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[GKP14]

[Proof in progress with Gourevitch, Kleinschmidt, Persson, Sahi]

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Would allow us to compute non-perturbative effects that capture information about instantons and black holes

[GKP14]

[Proof in progress with Gourevitch, Kleinschmidt, Persson, Sahi]

Outlook

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- Other compactifications leading to automorphic forms on other groups. (more conjectural)

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[Fleig-Kleinschmidt, Fleig-Kleinschmidt-Persson]

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How to define “small automorphic representations” for Kac-Moody groups? What is the mechanism behind the vanishing properties?

- $\mathcal{E}_6 D^6 R^4$ requires extended notion of automorphic forms, the development of which will positively bring new exciting insights to both physics and mathematics.

Thank you!

Henrik Gustafsson

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